Quantum Boltzmann Machines

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Can we encode complicated models into a relatively small neural network?

- quantum algorithm
- exploiting the full power of a quantum computer

• can handle quantum input/output

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• can handle quantum input/output

Boltzmann machines are a good candidate for quantizations [M. H. Amin, et al., arXiv:1601.02036]

What is a quantum Boltzmann machine?

Boltzmann machines are neural networks used for generative machine learning. [G. E. Hinton, 2012]



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Generative Training

"What I cannot create, I do not understand."

-Richard Feynman

Generative Training





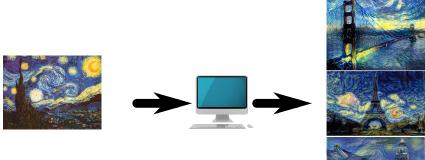
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[https://github.com/jcjohnson/neural-style]

Generative Training





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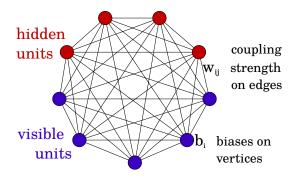
Algorithm learns a model that explains how the data were generated.

Assume there is an underlaying distribution of the data parametrized by {Θ_i}

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- Find $\{\Theta_i\}$ that "explains" the data
- Generate similar examples using the model

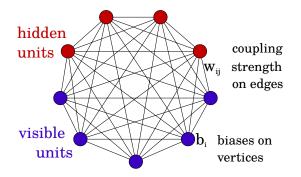
The Neural Network



visible units serve as input/output hidden units provide extra degrees of freedom

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The Neural Network



weights and biases are learned during training

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Boltzmann Machine

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Units s_i can take values -1 and 1Energy of a configuration

$$E(v,h) = \sum_{\text{vertices } i} b_i s_i + \sum_{\text{edges } \langle i,j \rangle} w_{i,j} s_i s_j$$

Motivated by the Ising model

$$E(v,h) = (-1)\cdot 1 + (-2)\cdot 1 + 3\cdot (-1) + + 3\cdot 1\cdot 1 + (-2)\cdot 1\cdot (-1) + 5\cdot 1\cdot (-1) = -6$$

Probability of a configuration on visible (v)and hidden (h) units

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$$p(v,h) = \frac{1}{Z}e^{-E(v,h)}$$
$$Z = \sum_{v,h} e^{-E(v,h)}$$

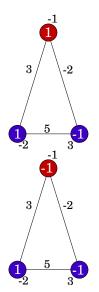
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Boltzmann distribution favors low-energy states

Probability of a configuration on visible (v)units is the marginal distribution



$$p(v) = \sum_{h} p(v, h) = \sum_{h} \frac{1}{Z} e^{-E(v, h)}$$
$$Z = \sum_{v, h} e^{-E(v, h)}$$

The distribution p(v) should be close to the distribution over the data q(v) for a properly trained Boltzmann machine.

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"Distance" Measure

The distribution $p(v) = \sum_{h} p(v, h)$ should be **close** to the distribution over the data q(v) for a properly trained Boltzmann machine **KL divergence**

$$\mathcal{L_{KL}} = \sum_{v \in data} q(v) \left[\log \left(q(v) \right) - \log \left(\sum_{h} p(v,h) \right) \right]$$

Negative log-likelihood

$$\mathcal{L} = -\sum_{v \in data} q(v) \log \left(\sum_{h} p(v, h)\right)$$

 \mathcal{L} is difficult to compute



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Minimize KL divergence using gradient descent

Gradient of \mathcal{L} is easy* to compute Requires only expectation values of single vertices or edges



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$$\frac{\partial \log p(v)}{\partial w_{i,j}} = \langle s_i s_j \rangle_{data} - \langle s_i s_j \rangle_{mode}$$
$$\frac{\partial \log p(v)}{\partial b_i} = \langle s_i \rangle_{data} - \langle s_i \rangle_{model}$$

*Given an approximation of a thermal state

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- 1. decide on the graph
- 2. generate starting weights and biases
- 3. construct the energy function
- 4. **repeat** for number of epochs:

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 - 4.1 measure expectation values in the thermal state

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- 4.2 repeat for each training example
 - 4.2.1 set visible units to the example vector
 - 4.2.2 create a thermal state on hidden units
 - 4.2.3 measure expectation values

- 1. decide on the graph
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- 4.2 repeat for each training example
 - 4.2.1 set visible units to the example vector
 - 4.2.2 create a thermal state on hidden units
 - 4.2.3 measure expectation values
- 4.3 compute gradients
- 4.4 update weights and biases
- 4.5 construct the new energy function

Quantum Boltzmann Machine

Quantum version of:

Energy function Objective function Training data

Energy Function

Classical thermal distribution:

$$p = \frac{1}{Z}e^{-E}, \quad Z = \sum_{z}e^{-E}$$

Quantum thermal state:

$$\rho = \frac{1}{Z}e^{-H}, \quad Z = Tr\left[e^{-H}\right]$$

Hamiltonian for Ising model

$$H = \sum_{i} b_i \sigma_i^z + \sum_{i < j} w_{i,j} \sigma_i^z \sigma_j^z$$

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energy $H \left| \psi \right\rangle = E \left| \psi \right\rangle$

Hamiltonian

Arbitrary QBM Hamiltonian

$$H = \sum_{i} \theta_i H_i, \qquad \|H_i\| = 1$$

Transverse Ising Model Pauli Basis

Fermionic Hamiltonian

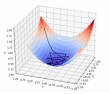
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Objective Function



KL divergence

$$\mathcal{L_{KL}} = \sum_{v \in data} q(v) \left[\log \left(q(v) \right) - \log \left(\sum_{h} p(v,h) \right) \right]$$

Quantum Relative Entropy

$$S = \operatorname{Tr} \left[\rho_{data} \left(\log \rho_{data} - \log \sigma_{model} \right) \right]$$

Objective Function

Minimize the objective function

$$\mathcal{O}_{\rho} = -\mathrm{Tr}\left[\rho_{data}\log\rho_{model}\right]$$

For no hidden units:

$$\partial_{\theta} \mathcal{O}_{\rho} = \langle \partial_{\theta} H \rangle_{qbm} - \langle \partial_{\theta} H \rangle_{data}$$

Learning Quantum States

$$\langle \partial_{\theta} H \rangle_{model} = \operatorname{Tr} \left[\partial_{\theta} H \frac{e^{-H}}{Tr[e^{-H}]} \right]$$

estimated by sampling from the QBM

$$\langle \partial_{\theta} H \rangle_{data} = Tr[\partial_{\theta} H \rho_{data}]$$

created by simulation



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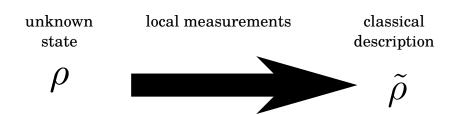
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Training data is a density matrix

Creating a representation of a quantum state is the goal of tomography

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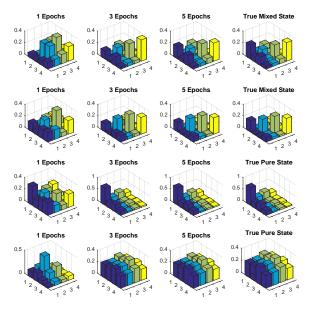
Tomography



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Tomography

- local measurements give gradients of relative entropy
- QBM describes the unknown state as a thermal state of a Hamiltonian
- Hamiltonian gives a classical approximate description of a state
- bonus: QBM works as an approximate cloning device



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- low number of epochs
- "compact" models for quantum states
- general method but can be tailored
- ability to create copies of a quantum state



- low number of epochs
- "compact" models for quantum states
- general method but can be tailored
- ability to create copies of a quantum state



- each epoch requires estimation of several expectations values
- approximating a thermal state can be costly

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- learn classical or quantum data
- gradient descent algorithm
- unorthodox approach to tomography
- approximate cloning

Objective Function - Part Deux

log-likelihood

$$\mathcal{L} = \sum_{v \in data} q(v) \log \left(\frac{\sum_{h} e^{-E(v,h)}}{\sum_{v',h'} e^{-E(v',h')}} \right)$$

quantum log-likelihood

$$\mathcal{O}_{\Lambda}(H) = \sum_{\mathbf{v}} P_{\mathbf{v}} \log \left(\frac{\operatorname{Tr} \left[\Lambda_{v} e^{-H} \right]}{\operatorname{Tr} \left[e^{-H} \right]} \right)$$

where Λ_v is a projector on state v on visible units

[M. H. Amin, et al., arXiv:1601.02036]



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Golden-Thompson Training

The gradient cannot be directly computed but can be bounded using Golden-Thomson inequality

$$\mathcal{L} \ge \sum_{v} P_{v} \log \left(\frac{Tr[\Lambda_{v}e^{-H_{v}}]}{Tr[e^{-H}]} \right)$$

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"clamped" Hamiltonian $H_v = H - ln\Lambda_v$ cannot learn states where $\text{Tr}[e^{-H_v}\partial_{\theta}H] = 0$ [M. H. Amin, et al., arXiv:1601.02036] How to extend the training set beyond classical states?

The set $\{\Lambda_v\}$ must be POVM The choice for classical data is ambiguous



Example: Training Set

• classical

$$\Lambda_n = |2n\rangle \langle 2n| \text{ for } 1 \le n \le 8$$

$$\Lambda_0 = \mathbb{1} - \sum_{n=1}^8 \Lambda_n, \qquad P_v = (1 - \delta_{v,0})/8.$$

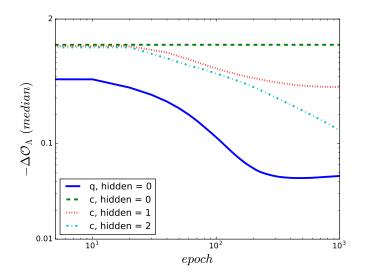
• superposition

$$\Lambda_1 = \frac{1}{8} (|2\rangle + \dots + |16\rangle) (\langle 2| + \dots + \langle 16|),$$

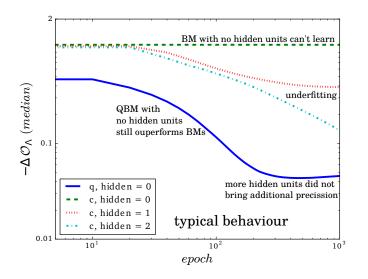
$$\Lambda_0 = \mathbb{1} - \Lambda_1, \qquad P_v = \delta_{v,1}.$$

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5 visible units



5 visible units



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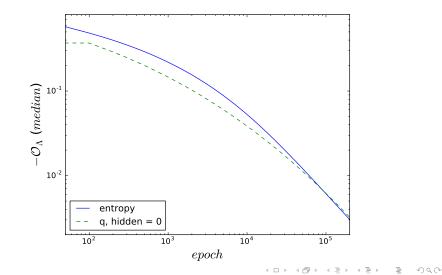
Golden-Thompson Training

- natural generalization of log-likelihood
- bound on the gradient can provide high-precision results

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• compact models of data

Quantum Relative Entropy vs Golden-Thompson Training



The thermal state is difficult to compute

Boltzmann machines face the same limitations but work well with a weak approximation of the thermal state (contrastive divergence for RBM)



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Thermal State

In practice we do not compute expectations values in the thermal state

- finite number of samples *n*
- approximation of the thermal state error ϵ_H

Error in estimating each component of the gradient

$$\mathcal{O}\left(\sqrt{1/n+\epsilon_H^2}\right)$$

Existing algorithms:

[A. N. Chowdhury and R. D. Somma, arXiv:1603.02940] [M. Yung and A. Aspuru-Guzik, Proc. Natl. Acad. Sci]

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Complexity

QBM is BQP hard

- can't be simulated efficiently unless BPP=BQP
- likely to provide quantum advantage

To Do:

- larger scale needed to show performance in practice
- unclear how approximations affect the convergence in practice

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Summary

- Potential application for medium size quantum computers
- Richer models, better approximation than classical Boltzmann machines

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• A new type of tomography - thermal state representation



- Potential application for medium size quantum computers
- Richer models, better approximation than classical Boltzmann machines
- A new type of tomography thermal state representation

